

EXAM QUANTUM THEORY, 24 JANUARY 2022, 14.15-17.15 HOURS.

1. In class we derived Kramers theorem for a Hermitian operator. Here we will derive a similar theorem for a unitary operator  $U$ . Let  $\mathcal{T}$  be an antiunitary operator that squares to  $-1$ . Assume that

$$\mathcal{T}U = U^{-1}\mathcal{T}.$$

- a) Let  $\Psi$  be an eigenstate of  $U$  with eigenvalue  $\lambda$ . Prove that  $\Psi' = \mathcal{T}\Psi$  is also an eigenstate of  $U$  with the same eigenvalue.
  - b) Prove that  $\Psi$  and  $\Psi'$  are *not* linearly related, so that indeed the eigenvalue  $\lambda$  is doubly degenerate.
2. In class we encountered coherent states, which are eigenstates of the (bosonic) annihilation operator  $a$ . Alice and Bob wonder about eigenstates of the creation operator  $a^\dagger$ . Bob says: If  $|\Psi\rangle$  is an eigenstate of  $a$  with eigenvalue  $\lambda$ , then  $|\Psi\rangle^*$  is an eigenstate of  $a^\dagger$  with eigenvalue  $\lambda^*$ . Alice disagrees.
- a) Who is right, Alice or Bob? Motivate your answer.
  - b) Suppose that  $|\beta\rangle$  is an eigenstate of  $a^\dagger$  with eigenvalue  $\beta \neq 0$ . Prove that  $\langle n+1|\beta\rangle = 0$  if  $\langle n|\beta\rangle = 0$ , for any number state  $|n\rangle$ .
  - c) Prove that  $\langle 0|\beta\rangle = 0$ . What do you conclude about the existence of eigenstates of the creation operator? What about the special case  $\beta = 0$ ?
3. We study a harmonic oscillator (frequency  $\omega$ ) with a time dependent mass\*  $m(t) = m_0 e^{\nu t}$ , for some real constants  $m_0$  and  $\nu$ . Its Hamiltonian is

$$H(t) = \frac{p^2}{2m(t)} + \frac{1}{2}m(t)\omega^2 x^2.$$

(We work in the Schrödinger picture, so the operators  $x$  and  $p$  are not time dependent.)

- a) Define the operators  $A = \frac{1}{2}\nu px + \frac{1}{4}i\hbar\nu$  and  $U(t) = e^{iAt/\hbar}$ . Is  $A$  Hermitian? Is  $U(t)$  unitary?
- b) Define  $P(t) = U(t)pU^{-1}(t)$ ,  $X(t) = U(t)xU^{-1}(t)$ . Derive that

$$\frac{d}{dt}P(t) = -\frac{\nu}{2}P(t), \quad \frac{d}{dt}X(t) = \frac{\nu}{2}X(t).$$

*Hint: notice that  $[A, p] = \frac{1}{2}i\hbar\nu p$  and  $[A, x] = -\frac{1}{2}i\hbar\nu x$ .*

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\*This is known as the Caldirola-Kanai oscillator, and is used as a model for damping by friction.

- *c)* Solve these differential equations to obtain  $P(t)$  and  $X(t)$  in terms of  $p$  and  $x$ . Then prove that

$$H(t) = U(t)H(0)U^{-1}(t).$$

- *d)* Charlie looks at this expression for  $H(t)$  and concludes that the wave function  $\psi(t)$  of the harmonic oscillator depends on time as

$$\psi(t) = U(t) \exp[-(it/\hbar)H(0)]\psi(0).$$

This is not quite correct, what term has Charlie overlooked? Explain why Charlie's equation can be called the "adiabatic approximation" for  $\psi(t)$ .

4. A particle of mass  $m$  moves freely along the  $x$ -axis, with Hamiltonian  $H(x, p) = \frac{1}{2}p^2/m$  and Lagrangian  $L(x, \dot{x}) = \frac{1}{2}m\dot{x}^2$ .

- *a)* Calculate the classical action  $S_{\text{class}} = \int_{t_1}^{t_2} L dt$  for the classical path from point  $x_1$  at time  $t_1$  to point  $x_2$  at time  $t_2$ .
- *b)* Calculate the quantum mechanical propagator<sup>†</sup>

$$G(x_2, t_2; x_1, t_1) = \langle x_2 | e^{-(i/\hbar)(t_2-t_1)\hat{H}} | x_1 \rangle.$$

- *c)* Discuss the relation between  $G$  and  $S_{\text{class}}$  in the context of Feynman's path integral formula.

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<sup>†</sup>You may use the integral  $\int_{-\infty}^{\infty} e^{ias-ibs^2} ds = \sqrt{\frac{\pi}{ib}} \exp\left(\frac{ia^2}{4b}\right)$ .